

# Effects of Pertinent Parameters on Pressure and Derivative Behaviour of a Layered Reservoir System with Crossflow

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**Abstract:** This paper studies the effects of some pertinent well and reservoir parameters on dimensionless pressure ( $P_D$ ) and derivative ( $P_D'$ ) behaviour of horizontal wells in a two layered reservoir subject to a bottom water drive. Pressure and derivative values were generated using layered reservoir models with a permeable interface. Results showed the emergence of the typical constant pressure boundary trend (constant pressure and zero derivative values) at late times. It was observed that this characteristic phenomenon shielded the effect of the parameters examined at late times. Both layers displayed some similar pressure and derivative behaviour with the varied parameters. The dimensionless pay thickness ( $h_D$ ), dimensionless wellbore radius ( $r_{wD}$ ) and crossflow coefficient ( $\beta$ ) influenced the pressure behaviour of the wells while the modification factor ( $E$ ), well width ( $y_D$ ) flow point ( $x_D$ ), well location ( $z_{wD}$ ) showed negligible effects.

Results from this study will find use in efficient well planning and effective depletion strategy as well as enhancing the overall performance of layered crossflow reservoir system with bottom water drive.

**Keywords:** Permeable interface, layered reservoir, dimensionless pressure, pressure derivative, constant pressure boundary horizontal well

## 1. INTRODUCTION

Layered reservoirs usually consist of two groups; crossflow system where layers can hydro dynamically communicate at the contact places and no crossflow (commingled system) where communication is only via the wellbore [1-4]. Research has shown that crossflow reservoir system experience lower pressure than no crossflow system. And crossflow direction is always from areas of low permeability to that of higher permeability [5]. Researches done in layered crossflow system include development of well test analysis models, horizontal well technology [6-10], enhanced oil recovery and unconventional reservoir modelling [11-13]. Research on pressure behaviour of single and layered reservoir systems showed that they behave differently [14]. Some factors like well completion, reservoir parameters, rock and fluid properties play vital role in shaping their characteristics profile and impact on productivity [15,16]. This paper aims to investigate the effects of changes in well / reservoir parameters (radius, pay thickness, modification factor, well width, crossflow coefficient and well location) on layer pressure and derivative response for a cross flow layered reservoir system using horizontal wells.

## 2. METHODOLOGY / MODEL DESCRIPTION

This study extends the work done by Akpobi et al. [17]. They developed mathematical models of a two layered reservoir with bottom water drive to study the performance of horizontal wells placed in each layer as shown in Figure 1a. The Source and Greens function method was utilized to formulate the mathematical model equations. This work utilized the pressure and derivative models under the condition of crossflow interface to study the effects of parameters like radius, pay thickness, modification factor, well width, crossflow coefficient and well location. Summary of theory is in Appendix. Model equations were numerically solved [17,18]. Plots of dimensionless pressure and derivative versus dimensionless time were developed and used for the analysis.

Table1.Example Parameters

Parameter	Well Length L(ft)	Height h.(ft)	Well Standoff, Z <sub>w</sub> (ft)	Reservoir Length, X <sub>e</sub> (ft)	Reservoir Width, Y <sub>e</sub> (ft)	Elevation Allowance(ft)	Permeability Ratio K <sub>2</sub> :K <sub>1</sub> (md)
Value	1000	200	100	10000	6000	20	10

### 3. RESULTS /DISCUSSION

Data in Table1 was used in mathematical models( $r_w = 0.375\text{ft}$ ,  $C_t = 1.0\text{E-}6 \text{ psi}^{-1}$ ,  $\phi=0.2$ ,  $\mu=0.3\text{cp}$ ,  $\alpha =2$ , flow rate is assumed constant at  $q = 400 \text{ STB/day}$ ).Same data was used for layer1 and 2 computations. The reservoir layer is isotropic ( $k_x = k_y = k_z$ ),  $E=1$  and  $\beta=1$ [17]

#### 3.1. Effects of Dimensionless Pay Thickness ( $h_D$ )

Figure 1 provides the result of pressure and derivative response of changes in  $h_D$  in a crossflow system. Using varying values of  $h_D$  in the range of  $0.05 \leq h_D \leq 1$ . It was observed that pressure and derivative response increased with larger  $h_D$  values. At late times constant pressure was observed and derivative values collapsed to zero. Results showed similar trend for both layers of the system.

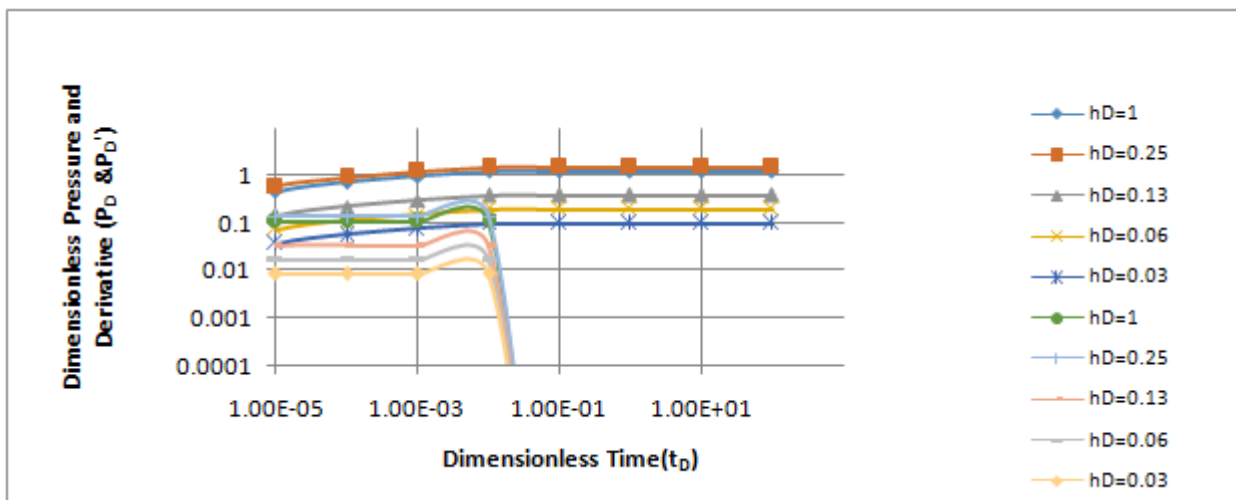


Figure1.Effects of dimensionless pay thickness ( $h_D$ )

#### 3.2. Influence of $Z_{WD}$

To study the influence of well location four values of  $z_{WD}$ (0.0, 0.125, 0.3 and 0.5) were used to compute pressure and derivative response. The single line for pressure and derivative in Figure 2 showed that pressure and derivative response were insensitive to changes in vertical well location.

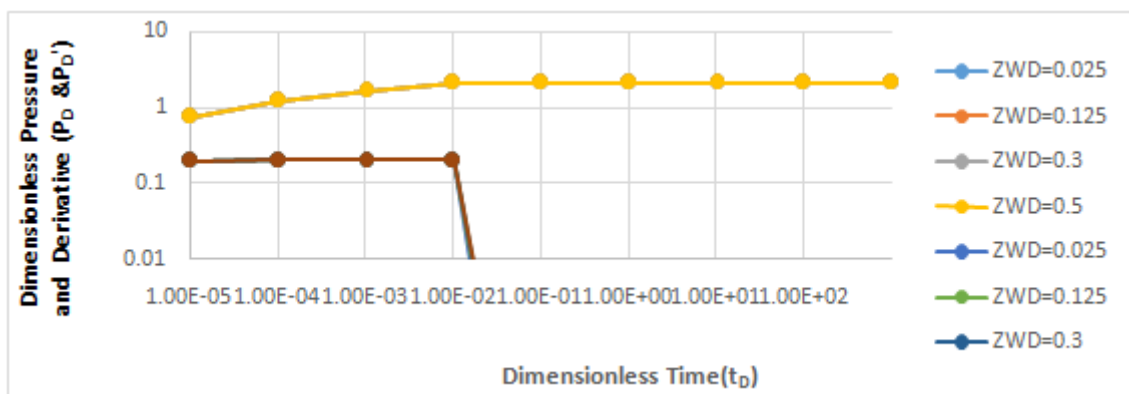


Figure2.Influence of  $Z_{WD}$

### 3.3. Influence of $r_{wD}$

Influence of changes in the wellbore radius on pressure and derivative response, was investigated using  $r_{wD}$  values in the range of  $1E-6 \leq r_{wD} \leq 1E-3$ . From Figure 3, it was observed that for larger  $r_{wD}$  pressure decreases, and this change was more pronounced during the early time flow period. Derivative was insensitive to changes in  $r_{wD}$ . Same results were obtained from layer two of the crossflow system. This result will find use in well completion for the system.

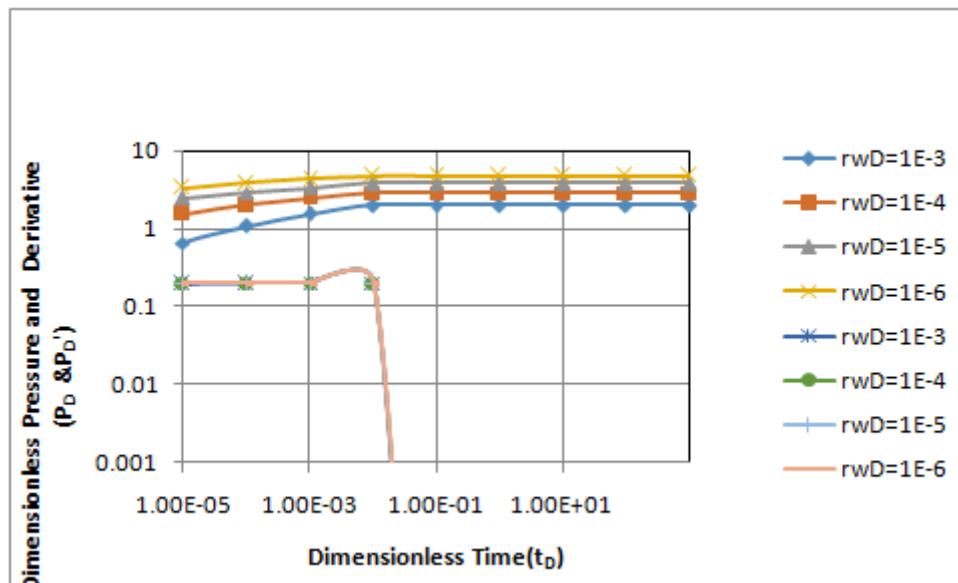


Figure3. Influence of  $R_{wD}$

### 3.4. Influence of $Y_D$

Influence of changes in the well width on pressure and derivative response, was investigated using  $y_D$  values in the range of  $5E-6 \leq y_D \leq 3E-3$ . As can be observed from Figure 4 changes in well width has no effect on pressure and derivative response for this cross flow layered system.

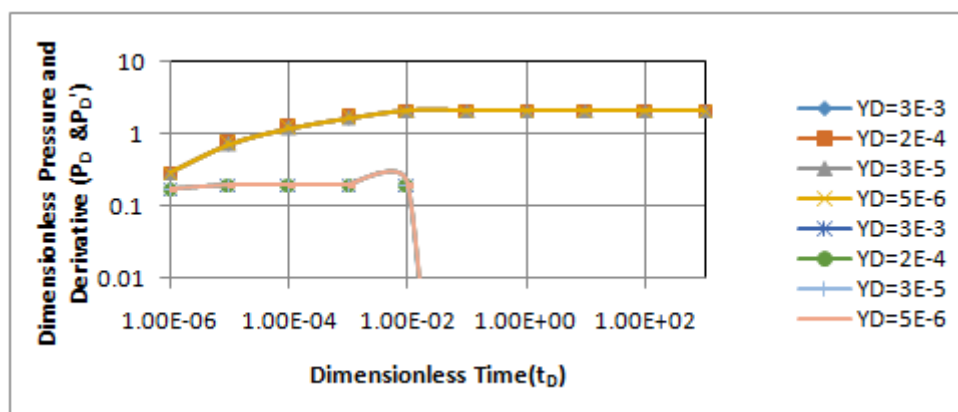


Figure4. Influence of  $Y_D$

### 3.5. Influence of $X_{wD}$

Also varying flow point ( $x_{wD}$ ) did not influence pressure and derivative as can be observed from Figure 5

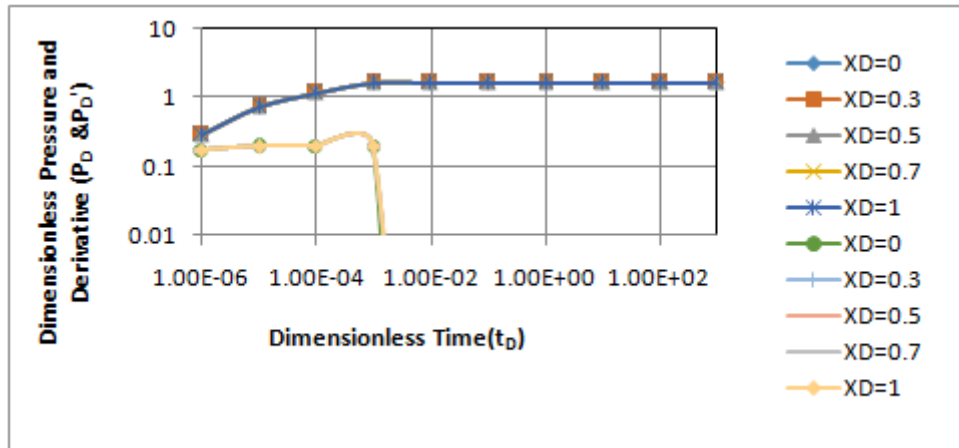


Figure 5. Influence of  $X_D$

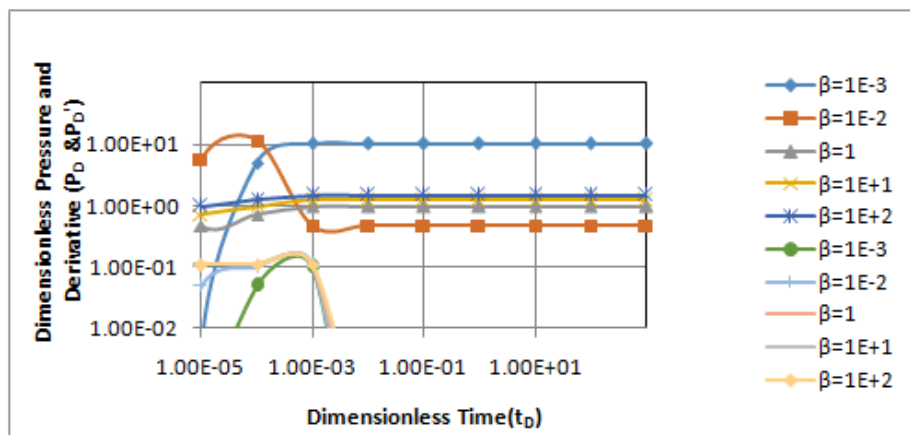


Figure 6. Influence of  $\beta$

### 3.6. Influence of $\beta$ (Crossflow Coefficient)

Values of  $1E-3 \leq \beta \leq 1E2$  were used to investigate its influence on pressure and derivative response and from figure 6, it was observed that for values of  $\beta \leq 1$  the early time pressure response was unstable but stable at late time. For  $\beta \geq 1$  there is little difference in pressure and derivative values although the values were stable. Hence the stable and favourable values of  $\beta$  are those greater than one (permeability of layer 2 must be higher). The crossflow direction is always from area of low permeability to that of higher permeability. Hence in the example, when there is drawdown, formation fluids start moving from layer with low  $k$  (layer 1) to that with a higher  $k$  (layer 2) as a result there is increased pressure response in layer two. Well should be completed in the more permeable layer.

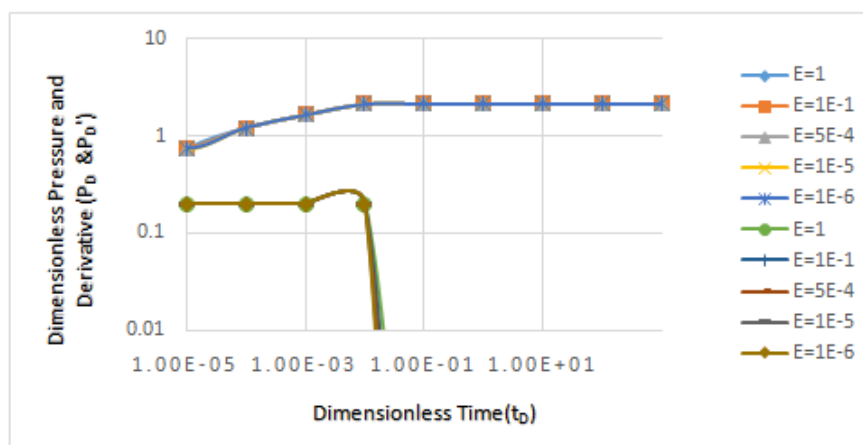


Figure 7. Influence of Modification Factor  $E$ .

### 3.7. Influence of Modification Factor E

Using values of  $1E-6 \leq E \leq 1$ . No effect was noticed on pressure and derivative response as a result of changes in the modification factor, as seen in Figure 7. E usually takes effect when early time flow period has elapsed, hence this may be attributed to constant pressure boundary effect. A late time trend of bottom water drive system where pressure stabilizes and derivatives collapses to zero. This observation was previously made by Ozkan [19] in his study of bottom water drive with horizontal wells for single systems.

## 4. CONCLUSION

Information on the effects of well and reservoir parameters on pressure and derivative behaviour can be utilized to improve performance of crossflow layered reservoir system. In this study the model equations for a layered crossflow reservoir system were presented. The effect of some pertinent well and reservoir parameters on the behaviour of pressure and pressure derivative of horizontal wells placed in each layer of the system were also examined. The following conclusions were made.

$P_D$  and  $P_D'$  response depicted the characteristic late time trend of a reservoir with constant pressure boundary. Using well one as the reference well,  $P_D$  and  $P_D'$  increased with increasing dimensionless pay thickness, while it decreased with increasing  $r_{WD}$ , more during the early time flow period. Early time response was unstable for values of  $\beta < 1$ , derivative was observed to show negligible changes.  $P_D$  and  $P_D'$  were insensitive to changes in  $Y_D$ ,  $X_D$  and E.

This study showed that pressure derivative is not a suitable tool to test for the sensitivity of these parameters. Both layers gave near similar results, but the study recommends that well should be completed in the more permeable layer of the system. Results from this study will find use in efficient well planning, completion and depletion strategy for layered system with bottom water drive.

## 5. NOMENCLATURE

- $C_t$  Total compressibility  $\text{Psi}^{-1}$
- h pay thickness Ft
- $h_{Di}$  total dimensionless pay thickness
- i distance in x,y or z direction
- K permeability md
- L well length ft
- M mobility ratio
- P pressure psi
- $P_D$  dimensionless pressure
- $P_D'$  dimensionless Pressure derivative
- psi pound per square inch
- $r_w$  well radius ft
- t time hrs
- $t_D$  dimensionless time
- $\beta$  crossflow coefficient
- $\phi$  porosity
- $\mu$  viscosity cp

## 6. SUBSCRIPT

- bt breakthrough
- D dimensionless
- e external
- t total
- w wellbore

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## APPENDIX

This section gives a brief summary of the theory and mathematical derivation involved in the model formulation as presented in the research done by Akpobi et al [9,17]. Dimensionless parameters used include dimensionless well and reservoir distances, time and pressure

$$i_D = \frac{2i}{L} \sqrt{\frac{k}{k_i}} \quad (1a)$$

(where  $i = x, y, z$ )

$$t_D = \frac{0.0010548Kt}{\phi\mu C_t L^2} \quad (2a)$$

$$P_D = \frac{kh\Delta p}{141.2q\mu B} \quad (3a)$$

The equation governing the flow of a slightly compressible fluid in a porous media is the diffusivity equation expressed in dimensionless form as

$$\frac{\partial^2 P_D}{\partial x^2} + \frac{\partial^2 P_D}{\partial y^2} + \frac{\partial^2 P_D}{\partial z^2} = \frac{\partial P_D}{\partial t_D} \quad (4a)$$

Fig 1a. Schematic of the model for a two layered Crossflow reservoir with horizontal wells

The pressure drop caused by production from a continuous source [20] is expressed as

$$\Delta P(x, y, z, t) = \frac{1}{\phi C_t} \int_0^t q_l \cdot S(x, y, z) \cdot \partial t \quad (5a)$$

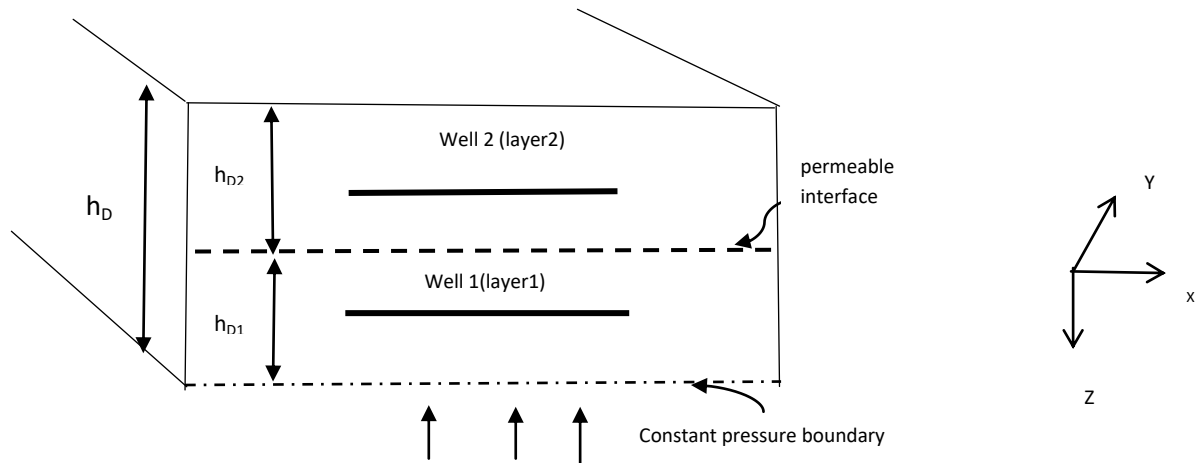


Figure 1a. Schematic of reservoir model

Where  $q_l$  represents flow rate per unit length of the source,  $s(x, y, z, t)$  represents the instantaneous source function (ISF) for the particular reservoir and well configuration [21]

$$P_D(X_D, Y_D, Z_D, \tau) = 2\pi h_D E_j \int_0^{\tau} S(X_D, \tau) \cdot S(Y_D, \tau) \cdot S(Z_D, \tau) \partial \tau \quad (6a)$$

Dimensionless Pressure ( $P_{Dj}$ ) for each layer is developed as

$$P_{Dj} = P_{D1} + P_{D2} + P_{D3} + \dots + P_{Dn} \quad (7a)$$

where  $j$  depicts the layer and  $n$  is the last flow period.

- 1. Modification Factor E** that accounts for the contributions of both layers at the interface as a result of the dual nature of the  $z$ -axis

$$EJ = \frac{W_1 + W_2}{2\pi(W_1^2 + W_2^2)} \quad (8a)$$

$$W_1 = h_{D1} \int_{X_{wD}}^{X_{eD}} \int_0^{y_{eD}} \int_0^{h_{D1}} \left[ \int_0^{t_D} s_1(x_D, \tau), s_1(y_D, \tau), s_1(z_D, \tau) d\tau \right] d_{zD} d_{yD} d_{xD} \quad (9a)$$

$$W_2 = h_{D2} \int_{X_{wD}}^{X_{eD}} \int_0^{y_{eD}} \int_{h_{D2}}^{h_D} \left[ \int_0^{t_D} s_2(x_D, \tau), s_2(y_D, \tau), s_2(z_D, \tau) d\tau \right] d_{zD} d_{yD} d_{xD} \quad (10a)$$



2. **Crossflow Coefficient ( $\beta$ );** is a factor used to correct for differences in response time of the layers to the same transient regime for crossflow systems.

$$\beta = \frac{\phi_i \mu_i C_{ii} L_i^2 K_{i+1}}{\phi_{i+1} \mu_{i+1} C_{i+1} L_{i+1}^2 K_i} \quad (11a)$$

3. **Total Permeability of the system**

$$K_t = \frac{K_1 h_1 + K_2 h_2 + \dots + K_n h_n}{h_1 h_2 + \dots + h_n} \quad (12a)$$

4. **Fluid mobility ratio**

$$M = \frac{K_1 h_1 / \mu_1}{K_2 h_2 / \mu_2} \quad (13a)$$

5. **Total pay thickness of the system**

$$h_{D1} = h_{D1} + h_{D2} \quad (14a)$$

Dimensionless pressure ( $P_{D1}$ ) for layer one crossflow systems

$$P_{D1} = \left[ \begin{aligned} & \left[ \frac{\alpha}{8L_D} \sqrt{\frac{K}{K_y}} Ei \left( -\frac{r_{WD}^2}{4\tau_D} \right) \right] + \\ & \left[ \left\{ \operatorname{erf} \left( \frac{\sqrt{\frac{K}{K_X}} + X_{D1}}{2\sqrt{\tau_D}} \right) + \operatorname{erf} \left( \frac{\sqrt{\frac{K}{K_X}} - X_{D1}}{2\sqrt{\tau_D}} \right) \right\} \right. \\ & \left. \sqrt{\pi} E_1 \int_{t_{De1}}^{t_{De2}} \left\{ 1 + 2 \sum_{n=1}^{\infty} \exp \left[ -\frac{n^2 \pi^2 \tau_D}{Y_{eD}^2} \right] \cos \frac{n\pi y_{wD}}{Y_{eD}} \cos \frac{n\pi Y_W}{Y_{eD}} \right\} \partial \tau + \right. \\ & \left. \frac{1}{2\sqrt{\tau_D}} \exp \left[ -\frac{(Z_{D1} - Z_{WD1})^2}{4\tau_D} \right] \right] \\ & \left[ \left\{ \operatorname{erf} \left( \frac{\sqrt{\frac{K}{K_X}} + X_{D1}}{2\sqrt{\tau_D}} \right) + \operatorname{erf} \left( \frac{\sqrt{\frac{K}{K_X}} - X_{D1}}{2\sqrt{\tau_D}} \right) \right\} \frac{1}{2\sqrt{\tau_D}} \exp \left[ -\frac{(Y_{D1} - Y_{WD1})^2}{4\tau_D} \right] \right. \\ & \left. \sqrt{\pi} E_1 \int_{t_{De1}}^{t_{De2}} \left\{ \sum_{n=1}^{\infty} \exp \left[ -\frac{n^2 \pi^2 \tau_D}{h_{D1}^2} \right] \sin \frac{n\pi Z_{wD1}}{h_{D1}} \sin \frac{n\pi z_{D1}}{h_{D1}} \right\} \partial \tau + \right. \\ & \left[ \left\{ 1 + \frac{4X_{eD}}{\pi} \sum_{m=1}^{\infty} \exp \left[ -\frac{m^2 \pi^2 \tau_D}{X_{eD}^2} \right] \sin \frac{m\pi x_f}{2X_{eD}} \cos \frac{m\pi x_{wD1}}{X_{eD}} \cos \frac{m\pi x_{D1}}{X_{eD}} \right\} \right. \\ & \left. \frac{4\pi E_1}{X_{eD} Y_{eD}} \int_{t_{De3}}^{t_D} \left\{ 1 + 2 \sum_{n=1}^{\infty} \exp \left[ -\frac{n^2 \pi^2 \tau_D}{Y_{eD}^2} \right] \cos \frac{n\pi y_{wD}}{Y_{eD}} \cos \frac{n\pi Y_W}{Y_{eD}} \right\} \partial \tau \right. \\ & \left. \left\{ \sum_{n=1}^{\infty} \exp \left[ -\frac{n^2 \pi^2 \tau_D}{h_{D1}^2} \right] \sin \frac{n\pi Z_{wD1}}{h_{D1}} \sin \frac{n\pi z_{D1}}{h_{D1}} \right\} \right] \end{aligned} \right] \quad (15a)$$







Dimensionless pressure derivative for layer two (crossflow systems) is given by the expression:

$$P_{D2}' = \left[ \begin{aligned} & \left[ \frac{\alpha}{8L_D} \sqrt{\frac{K}{K_y} \frac{K}{K_z}} \frac{1}{\tau_D} \exp\left(-\frac{r_{WD}^2}{4\tau_D}\right) + \right. \\ & \left. \sqrt{\pi} \beta E_2 \left\{ \left[ \operatorname{erf}\left(\frac{\sqrt{\frac{K}{K_x}} + X_{D2}}{2\sqrt{\beta\tau_D}}\right) + \operatorname{erf}\left(\frac{\sqrt{\frac{K}{K_x}} - X_{D2}}{2\sqrt{\beta\tau_D}}\right) \right] \right. \right. \\ & \left. \left. + \frac{1}{2\sqrt{\beta\tau_D}} \exp\left(-\frac{(Z_{D2}-Z_{WD2})^2}{4\beta\tau_D}\right) \left[ 1 + 2 \sum_{n=1}^{\infty} \exp\left(-\frac{n^2\pi^2\beta\tau_D}{Y_{eD}^2}\right) \cos\frac{n\pi y_{wD2}}{Y_{eD}} \cos\frac{n\pi Y_{D2}}{Y_{eD}} \right] \right\} + \right. \\ & \left. \sqrt{\pi} \beta E_2 \left\{ \left[ \operatorname{erf}\left(\frac{\sqrt{\frac{K}{K_x}} + X_{D2}}{2\sqrt{\beta\tau_D}}\right) + \operatorname{erf}\left(\frac{\sqrt{\frac{K}{K_x}} - X_{D2}}{2\sqrt{\beta\tau_D}}\right) \right] \cdot \frac{1}{2\sqrt{\beta\tau_D}} \exp\left(-\frac{(Y_{D2}-Y_{WD2})^2}{4\beta\tau_D}\right) \right. \right. \\ & \left. \left. + \sum_{n=1}^{\infty} \exp\left[-\frac{(2n-1)^2\pi^2\beta\tau_D}{4h_{D2}^2}\right] \sin\frac{(2n-1)\pi z_{wD2}}{2h_{D2}} \sin\frac{(2n-1)\pi z_{D2}}{2h_{D2}} \right] \right\} + \right. \\ & \left. \frac{4\pi\beta E_2}{X_{eD}Y_{eD}} \left\{ \left[ 1 + \frac{4X_{eD}}{\pi} \sum_{m=1}^{\infty} \exp\left[-\frac{m^2\pi^2\beta\tau_D}{X_{eD}^2}\right] \sin\frac{m\pi x_f}{2X_{eD}} \cos\frac{m\pi x_{wD2}}{X_{eD}} \cos\frac{m\pi x_{D2}}{X_{eD}} \right] \right. \right. \\ & \left. \left. + \frac{1}{2\sqrt{\beta\tau_D}} \exp\left(-\frac{(Z_{D2}-Z_{WD2})^2}{4\beta\tau_D}\right) \left[ 1 + 2 \sum_{n=1}^{\infty} \exp\left(-\frac{n^2\pi^2\beta\tau_D}{Y_{eD}^2}\right) \cos\frac{n\pi y_{wD2}}{Y_{eD}} \cos\frac{n\pi Y_{D2}}{Y_{eD}} \right] \right\} \right. \\ & \left. \left. + \sum_{n=1}^{\infty} \exp\left[-\frac{(2n-1)^2\pi^2\beta\tau_D}{4h_{D2}^2}\right] \sin\frac{(2n-1)\pi z_{wD2}}{2h_{D2}} \sin\frac{(2n-1)\pi z_{D2}}{2h_{D2}} \right] \right\} \right] \quad (18a) \end{aligned}$$

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